

A NOTE ON THE TENSOR FORM OF CHARACTERISTIC EQUATION OF 4×4 MATRIX USED IN HARISH-CHANDRA'S PAPER 'ALGEBRA OF DIRAC-MATRICES'

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Abstract

As the benefits of his beautiful formula for the product of Dirac-matrices Prof. Harish-Chandra evaluated the characteristic equation of a 4×4 matrix in his paper [1]. A trial to understand his ideas and his methods of calculation gives an internal pleasure at least to me. Through this paper I want to convey that pleasure to others also. The analysis given in this paper is useful in studying the relativistic wave equation of an electron.

Keywords: Characteristic polynomial, Characteristic equation, Trace, Dirac-matrices, Antisymmetric tensor.

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1. INTRODUCTION:

Let T be any 4×4 matrix. The matrix, $(T - tI)$ where I is the 4×4 identity matrix and t is an indeterminate, is called the characteristic matrix of T . Its determinant $\Delta(t) = \det(T - tI)$, which is a polynomial in t , is called the characteristic polynomial of T and $\Delta(t) = \det(T - tI) = 0$ is called the characteristic equation of T . The general characteristic polynomial of any 4×4 matrix T is of the form

$$\Delta(t) = t^4 - a_1 t^3 + a_2 t^2 - a_3 t + a_4, \quad (1)$$

where a_1, a_2, a_3, a_4 are constants s.t. $a_1 = \text{trace of } T = \text{tr}(T)$, $a_4 = \text{determinant value of } T = \det(T)$ and a_2, a_3 are the sum of the principal minors of order 2, 3 respectively.

If we put $t = 0$ in equation (1) we get the constant a_4 i.e. $\det(T)$ and also if we differentiate equation (1), three times with respect to t , then put $t = 0$ and divide by -6 , we get the constant a_1 i.e. $\text{tr}(T)$.

Harish-Chandra Formula [1] for the product of Dirac-matrices can be given as

$$E_{\lambda\mu} E_{\alpha\beta} = -\delta_{\lambda\alpha} \delta_{\mu\beta} + \delta_{\mu\alpha} \delta_{\lambda\beta} + E_{\lambda\alpha} \delta_{\mu\beta} - E_{\mu\alpha} \delta_{\lambda\beta} - E_{\lambda\beta} \delta_{\mu\alpha} + E_{\mu\beta} \delta_{\lambda\alpha} - \frac{i}{2} \varepsilon_{\lambda\mu\alpha\beta\gamma\rho} E^{\gamma\rho}, \quad (2)$$

where $E_{\lambda\mu}$ are Dirac matrices, $\delta_{\lambda\alpha}$ is a Kronecker delta function define as

$$\delta_{\lambda\alpha} = \begin{cases} 1, & \lambda = \alpha \\ 0, & \lambda \neq \alpha \end{cases}, \quad (3)$$

and $\varepsilon_{\lambda\mu\alpha\beta\gamma\rho}$ is an antisymmetric tensor in all six indices with $\varepsilon_{012345} = 1$. All the indices

$\lambda, \mu, \nu, \alpha, \beta, \gamma, \rho$ vary from 0 to 5. If two or more than two indices value is same the quantity $\varepsilon_{\lambda\mu\alpha\beta\gamma\rho}$ vanishes. Following paper [1] (2) is an equation in mixed tensor of rank four in a six dimensional space whose metric tensor is $\delta_{\lambda\alpha}$ for lowering or raising the index. The summation convention appear in a term if the same index appearing once above and below in the same term. For example

$$(i) E_{\lambda}^{\mu} E_{\mu\alpha} = \sum_{\mu=0}^5 E_{\lambda\mu} E_{\mu\alpha},$$

$$(ii) E^{\lambda\mu} \varepsilon_{\lambda\mu\alpha\beta\gamma\rho} = \sum_{\lambda, \mu=0}^5 E_{\lambda\mu} \varepsilon_{\lambda\mu\alpha\beta\gamma\rho},$$

(iii) but no summation is intended in term $E_{\lambda\mu} E_{\mu\alpha}$.

Dirac-matrices [1,2] $\{\gamma_1, \gamma_2, \gamma_3, \gamma_4\}$ are the set of four matrices of order 4×4 whose entries are complex numbers, given by the relation

$$\gamma_{\lambda} \gamma_{\mu} + \gamma_{\mu} \gamma_{\lambda} = 2\delta_{\lambda\mu}, \quad (\lambda, \mu = 1, 2, 3, 4) \quad (4)$$

these matrices are anticommute to each others with each square is identity matrix and generate a set of sixteen independent quantities which are also anticommute to each others (excluding identity matrix) can be given as

$$1; \gamma_1, \gamma_2, \gamma_3, \gamma_4; \gamma_1\gamma_2, \gamma_1\gamma_3, \gamma_1\gamma_4, \gamma_2\gamma_3, \gamma_2\gamma_4, \gamma_3\gamma_4; \gamma_1\gamma_2\gamma_3, \gamma_1\gamma_2\gamma_4, \gamma_1\gamma_3\gamma_4, \gamma_2\gamma_3\gamma_4; \gamma_1\gamma_2\gamma_3\gamma_4.$$

Since square of Dirac-matrices is 1 (identity matrix) but here some quantities square is -1 . So if we multiply i by these quantities then all 15 quantities satisfy the equation (4) and we call all of them Dirac-matrices and which can be given as

$$1; \gamma_1, \gamma_2, \gamma_3, \gamma_4; i\gamma_1\gamma_2, i\gamma_1\gamma_3, i\gamma_1\gamma_4, i\gamma_2\gamma_3, i\gamma_2\gamma_4, i\gamma_3\gamma_4; i\gamma_1\gamma_2\gamma_3, i\gamma_1\gamma_2\gamma_4, i\gamma_1\gamma_3\gamma_4, i\gamma_2\gamma_3\gamma_4; \gamma_1\gamma_2\gamma_3\gamma_4.$$

For simplicity and quantities which can be used in equation (2), we can change above matrices in one index notation, suppose

$$E_1 = i\gamma_1, E_2 = i\gamma_2, E_3 = i\gamma_3, E_4 = i\gamma_4,$$

$$E_5 = i(i\gamma_1\gamma_2), E_6 = i(i\gamma_1\gamma_3), E_7 = i(i\gamma_1\gamma_4), E_8 = i(i\gamma_2\gamma_3), E_9 = i(i\gamma_2\gamma_4), E_{10} = i(i\gamma_3\gamma_4),$$

$$E_{11} = i(i\gamma_1\gamma_2\gamma_3), E_{12} = i(i\gamma_1\gamma_2\gamma_4), E_{13} = i(i\gamma_1\gamma_3\gamma_4), E_{14} = i(i\gamma_2\gamma_3\gamma_4),$$

$$E_{15} = i(\gamma_1\gamma_2\gamma_3\gamma_4), E_{16} = i(1).$$

So equation (4) can be written as

$$E_\lambda E_\mu + E_\mu E_\lambda = -2\delta_{\lambda\mu}. \quad (\lambda, \mu = 1 \text{ to } 15,) \quad (5)$$

Hence i, E_λ ($\lambda = 1, 2, 3, \dots, 15$) form a complete set with properties $E_\lambda^2 = -1$ ($\lambda = 1$ to 16) and

$E_\lambda E_\mu = -E_\mu E_\lambda$, ($\lambda, \mu = 1$ to 15, $\lambda \neq \mu$). As in paper [1] we can change one index notation of these

sixteen quantities into two indices notation with help of following rules

$$E_{\lambda\mu} = E_\lambda E_\mu, \quad (\lambda \neq \mu) \quad (\lambda = 0, 1, 2, 3, 4, 5)$$

$$E_{\lambda\mu} = -E_{\mu\lambda}, \quad (\lambda \neq \mu) \quad (\lambda = 0, 1, 2, 3, 4, 5)$$

$$E_\lambda = E_{0\lambda} = -E_{\lambda 0}, \quad (\lambda = 1, 2, 3, 4, 5)$$

$$E_{\lambda\lambda} = 0. \quad (\lambda = 0, 1, 2, 3, 4, 5)$$

Then $i, E_{\lambda\mu}$ ($\lambda \neq \mu, \lambda, \mu = 0, 1, 2, 3, 4, 5$) form a complete set of Dirac-matrices with regard that two

quantities which differ by a numerical factor $-1, i$ or $-i$ is essentially the same i.e.

$$i, E_{\lambda\mu} = \left\{ \begin{array}{cccccc} E_{00} & E_{01} & E_{02} & E_{03} & E_{04} & E_{05} \\ E_{10} & E_{11} & E_{12} & E_{13} & E_{14} & E_{15} \\ E_{20} & E_{21} & E_{22} & E_{23} & E_{24} & E_{25} \\ E_{30} & E_{31} & E_{32} & E_{33} & E_{34} & E_{35} \\ E_{40} & E_{41} & E_{42} & E_{43} & E_{44} & E_{45} \\ E_{50} & E_{51} & E_{52} & E_{53} & E_{54} & E_{55} \end{array} \right\},$$

here the diagonal entries of $E_{\lambda\mu}$ are zero and entries below the diagonal are negative times corresponding entries in above the diagonal. So $i, E_{\lambda\mu}$ are consisting only sixteen distinct quantities $i, E_{01}, E_{02}, E_{03}, E_{04}, E_{05}; E_{12}, E_{13}, E_{14}, E_{15}; E_{23}, E_{24}, E_{25}; E_{34}, E_{35}; E_{45}$, which are anticommute to each others (excluding identity matrix i) and each square is -1 , form a complete set of matrices of order 4×4 . Hence these 16 matrices form a basis of a vector space of dimension 16 over the field of complex numbers. So any 4×4 matrix can be written as linear combination of these 16 matrices. This linear combination has been written in condensed form in next section with help of tensor calculus [3].

Prof. Harish-Chandra evaluated characteristic equation of any 4×4 matrix by using formula (2) with the help of two tensor identities [1]. The proof of these tensor identities have not given in his paper. I have verified these tensor identities and trying to prove them. These tensor identities can be given as

$$\epsilon_{\alpha\beta\gamma\delta\lambda\rho} \epsilon^{\lambda\rho\alpha'\beta'\gamma'\delta'} t_{\alpha'\beta'} t_{\gamma'\delta'} = 16(t_{\alpha\beta} t_{\gamma\delta} - t_{\alpha\gamma} t_{\beta\delta} - t_{\alpha\delta} t_{\gamma\beta}), \quad (6)$$

$$t_{\alpha\lambda} t_{\beta\rho} \epsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\mu\nu} t_{\sigma\tau} = \frac{1}{6} t_{\lambda\rho} \epsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\alpha\beta} t_{\mu\nu} t_{\sigma\tau}. \quad (7)$$

2. DISCUSSION:

Let T be any 4×4 matrix can be written [1] as

$$T = t + t_{\lambda\mu} E^{\lambda\mu}, \quad (8)$$

where $t, t_{\lambda\mu}$ are ordinary numbers commute to all other quantities s.t. $t_{\lambda\mu} = -t_{\mu\lambda}$ and $E_{\lambda\mu}$ are 15 Dirac-matrices and t represent identity matrix which multiply by number t . Now here we have first evaluated the characteristic polynomial of T . Equation (8) can be written as

$$T - t = t_{\lambda\mu} E^{\lambda\mu}, \quad (9)$$

squaring both side of above equation we get

$$\begin{aligned} (T - t)^2 &= (t_{\lambda\mu} E^{\lambda\mu})^2, \\ &= t_{\lambda\mu} E^{\lambda\mu} t_{\alpha\beta} E^{\alpha\beta} = \sum_{\lambda, \mu, \alpha, \beta=0}^5 t_{\lambda\mu} t_{\alpha\beta} E_{\lambda\mu} E_{\alpha\beta} \quad (\text{by using summation convention}) \\ &= \sum_{\lambda, \mu, \alpha, \beta=0}^5 t_{\lambda\mu} t_{\alpha\beta} \left\{ -\delta_{\lambda\alpha} \delta_{\mu\beta} + \delta_{\mu\alpha} \delta_{\lambda\beta} + E_{\lambda\alpha} \delta_{\mu\beta} - E_{\mu\alpha} \delta_{\lambda\beta} - E_{\lambda\beta} \delta_{\mu\alpha} + E_{\mu\beta} \delta_{\lambda\alpha} - \frac{i}{2} \epsilon_{\lambda\mu\alpha\beta\gamma\rho} E^{\gamma\rho} \right\}, \end{aligned}$$

here we use well known Harish-Chandra Formula (2) for the product of Dirac-matrices.

$$\begin{aligned} \text{Now } (T - t)^2 &= - \sum_{\lambda, \mu=0}^5 t_{\lambda\mu} t_{\lambda\mu} + \sum_{\lambda, \mu=0}^5 t_{\lambda\mu} t_{\mu\lambda} + \sum_{\lambda, \mu, \alpha=0}^5 t_{\lambda\mu} t_{\alpha\mu} E_{\lambda\alpha} - \sum_{\lambda, \mu, \alpha=0}^5 t_{\lambda\mu} t_{\alpha\lambda} E_{\mu\alpha} \\ &\quad - \sum_{\lambda, \mu, \beta=0}^5 t_{\lambda\mu} t_{\mu\beta} E_{\lambda\beta} + \sum_{\lambda, \mu, \beta=0}^5 t_{\lambda\mu} t_{\lambda\beta} E_{\mu\beta} - \frac{i}{2} \sum_{\lambda, \mu, \alpha, \beta=0}^5 t_{\lambda\mu} t_{\alpha\beta} \epsilon_{\lambda\mu\alpha\beta\gamma\rho} E^{\gamma\rho}, \end{aligned}$$

by means of $t_{\mu\lambda} = -t_{\lambda\mu}$ in second term and opening summation in third, fourth, fifth, sixth term with using $E_{\lambda\mu} = -E_{\mu\lambda}$, these four term vanishes then so that

$$(T - t)^2 = -2 \sum_{\lambda, \mu=0}^5 t_{\lambda\mu} t_{\lambda\mu} - \frac{i}{2} \sum_{\lambda, \mu, \alpha, \beta=0}^5 t_{\lambda\mu} t_{\alpha\beta} \epsilon_{\lambda\mu\alpha\beta\gamma\rho} E^{\gamma\rho} = -2t_{\lambda\mu} t^{\lambda\mu} - \frac{i}{2} t^{\lambda\mu} t^{\alpha\beta} \epsilon_{\lambda\mu\alpha\beta\gamma\rho} E^{\gamma\rho},$$

for convenience change summation indices in the above equation then we above

$$(T - t)^2 + 2t_{\mu\nu} t^{\mu\nu} = -\frac{i}{2} t^{\alpha\beta} t^{\gamma\delta} \epsilon_{\alpha\beta\gamma\delta\lambda\rho} E^{\lambda\rho},$$

again squaring both side of above equation we get

$$\left\{ (T - t)^2 + 2t_{\mu\nu} t^{\mu\nu} \right\}^2 = \left\{ -\frac{i}{2} t^{\alpha\beta} t^{\gamma\delta} \epsilon_{\alpha\beta\gamma\delta\lambda\rho} E^{\lambda\rho} \right\}^2$$

$$\begin{aligned}
 &= \left\{ -\frac{i}{2} t^{\alpha\beta} t^{\gamma\delta} \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} E^{\lambda\rho} \right\} \left\{ -\frac{i}{2} t^{\alpha'\beta'} t^{\gamma'\delta'} \varepsilon_{\alpha'\beta'\gamma'\delta'\mu\nu} E^{\mu\nu} \right\} \\
 &= \sum_{\lambda,\rho,\mu,\nu=0}^5 \left\{ -\frac{1}{4} t^{\alpha\beta} t^{\gamma\delta} \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} t^{\alpha'\beta'} t^{\gamma'\delta'} \varepsilon_{\alpha'\beta'\gamma'\delta'\mu\nu} E_{\lambda\rho} E_{\mu\nu} \right\} \quad (\text{by using formula (2) in } E_{\lambda\rho} E_{\mu\nu}) \\
 &= \sum_{\lambda,\rho,\mu,\nu=0}^5 \left\{ -\frac{1}{4} t^{\alpha\beta} t^{\gamma\delta} \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} t^{\alpha'\beta'} t^{\gamma'\delta'} \varepsilon_{\alpha'\beta'\gamma'\delta'\mu\nu} \left\{ \begin{aligned} &-\delta_{\lambda\mu} \delta_{\rho\nu} + \delta_{\rho\mu} \delta_{\lambda\nu} + E_{\lambda\mu} \delta_{\rho\nu} - E_{\rho\mu} \delta_{\lambda\nu} \\ &-E_{\lambda\nu} \delta_{\rho\mu} + E_{\rho\nu} \delta_{\lambda\mu} - \frac{i}{2} \varepsilon_{\lambda\rho\mu\nu\sigma\tau} E^{\sigma\tau} \end{aligned} \right\} \right\} \\
 &= -\frac{1}{4} t^{\alpha\beta} t^{\gamma\delta} t^{\alpha'\beta'} t^{\gamma'\delta'} \left\{ -\sum_{\lambda,\rho=0}^5 \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon_{\alpha'\beta'\gamma'\delta'\lambda\rho} + \sum_{\lambda,\rho=0}^5 \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon_{\alpha'\beta'\gamma'\delta'\rho\lambda} + \sum_{\lambda,\rho,\mu=0}^5 \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon_{\alpha'\beta'\gamma'\delta'\mu\rho} E_{\lambda\mu} \right. \\
 &\quad - \sum_{\lambda,\rho,\mu=0}^5 \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon_{\alpha'\beta'\gamma'\delta'\mu\lambda} E_{\rho\mu} - \sum_{\lambda,\rho,\nu=0}^5 \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon_{\alpha'\beta'\gamma'\delta'\rho\nu} E_{\lambda\nu} + \sum_{\lambda,\rho,\nu=0}^5 \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon_{\alpha'\beta'\gamma'\delta'\lambda\nu} E_{\rho\nu} \\
 &\quad \left. - \frac{i}{2} \sum_{\lambda,\rho,\mu,\nu=0}^5 \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon_{\alpha'\beta'\gamma'\delta'\mu\nu} \varepsilon_{\lambda\rho\mu\nu\sigma\tau} E^{\sigma\tau} \right\},
 \end{aligned}$$

we used summation convention in 1st, 2nd, 7th terms and 3rd, 4th terms cancel out the terms 5th, 6th respectively in the above equation then

$$\begin{aligned}
 &\left\{ (T-t)^2 + 2t_{\mu\nu} t^{\mu\nu} \right\}^2 \\
 &= -\frac{1}{4} t^{\alpha\beta} t^{\gamma\delta} (-2) \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon^{\lambda\rho\alpha'\beta'\gamma'\delta'} t_{\alpha'\beta'} t_{\gamma'\delta'} + \frac{i}{8} \varepsilon^{\alpha\beta\gamma\delta\lambda\rho} t_{\alpha\beta} t_{\gamma\delta} \varepsilon^{\alpha'\beta'\gamma'\delta'\mu\nu} t_{\alpha'\beta'} t_{\gamma'\delta'} \varepsilon_{\lambda\rho\mu\nu\sigma\tau} E^{\sigma\tau} \\
 &= \frac{1}{2} t^{\alpha\beta} t^{\gamma\delta} \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon^{\lambda\rho\alpha'\beta'\gamma'\delta'} t_{\alpha'\beta'} t_{\gamma'\delta'} + \frac{i}{8} \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\mu\nu} t_{\sigma\tau} \varepsilon^{\gamma\delta\mu'\nu'\sigma'\tau'} t_{\mu'\nu'} t_{\sigma'\tau'} \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} E^{\lambda\rho},
 \end{aligned}$$

for convenience here we changed summation indices in second term of right side of equation then above equation can be written as

$$\begin{aligned}
 &\left\{ (T-t)^2 + 2t_{\mu\nu} t^{\mu\nu} \right\}^2 - \frac{1}{2} t^{\alpha\beta} t^{\gamma\delta} \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon^{\lambda\rho\alpha'\beta'\gamma'\delta'} t_{\alpha'\beta'} t_{\gamma'\delta'} \\
 &= \frac{i}{8} \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\mu\nu} t_{\sigma\tau} \varepsilon^{\gamma\delta\mu'\nu'\sigma'\tau'} t_{\mu'\nu'} t_{\sigma'\tau'} \varepsilon_{\alpha\beta\lambda\rho\gamma\delta} E^{\lambda\rho} \\
 &= \frac{i}{8} \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\mu\nu} t_{\sigma\tau} \left\{ \varepsilon_{\alpha\beta\lambda\rho\gamma\delta} \varepsilon^{\gamma\delta\mu'\nu'\sigma'\tau'} t_{\mu'\nu'} t_{\sigma'\tau'} \right\} E^{\lambda\rho} \quad (\text{using first tensor identity (6)}) \\
 &= \frac{i}{8} \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\mu\nu} t_{\sigma\tau} \left\{ 16 \left\{ t_{\alpha\beta} t_{\lambda\rho} - t_{\alpha\lambda} t_{\beta\rho} - t_{\alpha\rho} t_{\lambda\beta} \right\} \right\} E^{\lambda\rho}
 \end{aligned}$$

$$\begin{aligned}
 &= 2i\varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\mu\nu} t_{\sigma\tau} \left\{ t_{\alpha\beta} t_{\lambda\rho} E^{\lambda\rho} - 2t_{\alpha\lambda} t_{\beta\rho} E^{\lambda\rho} \right\} && \text{(by using equation (9))} \\
 &= 2i\varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\alpha\beta} t_{\mu\nu} t_{\sigma\tau} (T-t) - 4it_{\alpha\lambda} t_{\beta\rho} \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\mu\nu} t_{\sigma\tau} E^{\lambda\rho}
 \end{aligned}$$

$$\begin{aligned}
 \text{So } \left\{ (T-t)^2 + 2t_{\mu\nu} t^{\mu\nu} \right\}^2 - \frac{1}{2} t^{\alpha\beta} t^{\gamma\delta} \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon^{\lambda\rho\alpha'\beta'\gamma'\delta'} t_{\alpha'\beta'} t_{\gamma'\delta'} \\
 = 2i\varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\alpha\beta} t_{\mu\nu} t_{\sigma\tau} (T-t) - 4it_{\alpha\lambda} t_{\beta\rho} \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\mu\nu} t_{\sigma\tau} E^{\lambda\rho}. \quad (10)
 \end{aligned}$$

$$\begin{aligned}
 \text{But the quantity } &-4it_{\alpha\lambda} t_{\beta\rho} \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\mu\nu} t_{\sigma\tau} E^{\lambda\rho} \\
 &= -4i \left\{ \frac{1}{6} t_{\lambda\rho} \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\alpha\beta} t_{\mu\nu} t_{\sigma\tau} \right\} E^{\lambda\rho} && \text{(using second tensor identity (7))} \\
 &= -\frac{4}{6} i \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\alpha\beta} t_{\mu\nu} t_{\sigma\tau} t_{\lambda\rho} E^{\lambda\rho} && \text{(again using equation (9))} \\
 &= -\frac{2}{3} i \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\alpha\beta} t_{\mu\nu} t_{\sigma\tau} (T-t).
 \end{aligned}$$

$$\begin{aligned}
 \text{Also the quantity } &-\frac{1}{2} t^{\alpha\beta} t^{\gamma\delta} \varepsilon_{\alpha\beta\gamma\delta\lambda\rho} \varepsilon^{\lambda\rho\alpha'\beta'\gamma'\delta'} t_{\alpha'\beta'} t_{\gamma'\delta'} && \text{(again using first tensor identity (6))} \\
 &= -\frac{1}{2} t^{\alpha\beta} t^{\gamma\delta} \left\{ 16 \left(t_{\alpha\beta} t_{\gamma\delta} - t_{\alpha\gamma} t_{\beta\delta} - t_{\alpha\delta} t_{\gamma\beta} \right) \right\} \\
 &= -8t^{\alpha\beta} t^{\gamma\delta} \left(t_{\alpha\beta} t_{\gamma\delta} - t_{\alpha\gamma} t_{\beta\delta} - t_{\alpha\delta} t_{\gamma\beta} \right) \\
 &= -8 \left\{ \left(t^{\alpha\beta} t_{\alpha\beta} \right)^2 - t^{\alpha\beta} (-t^{\gamma\delta}) (-t_{\delta\alpha}) t_{\beta\gamma} - t^{\alpha\beta} t^{\gamma\delta} (-t_{\delta\alpha}) (-t_{\beta\gamma}) \right\} \\
 &= -8 \left\{ \left(t^{\alpha\beta} t_{\alpha\beta} \right)^2 - t^{\alpha\beta} t_{\beta\gamma} t^{\gamma\delta} t_{\delta\alpha} - t^{\alpha\beta} t_{\beta\gamma} t^{\gamma\delta} t_{\delta\alpha} \right\} \\
 &= -8 \left\{ \left(t^{\alpha\beta} t_{\alpha\beta} \right)^2 - 2t^{\alpha\beta} t_{\beta\gamma} t^{\gamma\delta} t_{\delta\alpha} \right\}.
 \end{aligned}$$

Substituting these values in equation (10) so that

$$\left\{ (T-t)^2 + 2t_{\mu\nu} t^{\mu\nu} \right\}^2 - 8 \left\{ \left(t^{\alpha\beta} t_{\alpha\beta} \right)^2 - 2t^{\alpha\beta} t_{\beta\gamma} t^{\gamma\delta} t_{\delta\alpha} \right\} = \frac{4}{3} i \varepsilon^{\alpha\beta\mu\nu\sigma\tau} t_{\alpha\beta} t_{\mu\nu} t_{\sigma\tau} (T-t),$$

for convenience again changed the summation indices in right side of equation as

$$\left\{ (T-t)^2 + 2t_{\mu\nu} t^{\mu\nu} \right\}^2 - 8 \left\{ \left(t^{\alpha\beta} t_{\alpha\beta} \right)^2 - 2t^{\alpha\beta} t_{\beta\gamma} t^{\gamma\delta} t_{\delta\alpha} \right\} = \frac{4}{3} it_{\alpha\beta} t_{\gamma\delta} t_{\lambda\rho} \varepsilon^{\alpha\beta\gamma\delta\lambda\rho} (T-t). \quad (11)$$

So Characteristic polynomial of matrix T can be

$$\Delta(T) = \left\{ (T-t)^2 + 2t_{\mu\nu}t^{\mu\nu} \right\}^2 - 8 \left\{ \left(t^{\alpha\beta}t_{\alpha\beta} \right)^2 - 2t^{\alpha\beta}t_{\beta\gamma}t^{\gamma\delta}t_{\delta\alpha} \right\} - \frac{4}{3}it_{\alpha\beta}t_{\gamma\delta}t_{\lambda\rho}\epsilon^{\alpha\beta\gamma\delta\lambda\rho} (T-t), \quad (12)$$

also Characteristic equation of matrix T is

$$\left\{ (T-t)^2 + 2t_{\mu\nu}t^{\mu\nu} \right\}^2 - 8 \left\{ \left(t^{\alpha\beta}t_{\alpha\beta} \right)^2 - 2t^{\alpha\beta}t_{\beta\gamma}t^{\gamma\delta}t_{\delta\alpha} \right\} - \frac{4}{3}it_{\alpha\beta}t_{\gamma\delta}t_{\lambda\rho}\epsilon^{\alpha\beta\gamma\delta\lambda\rho} (T-t) = 0, \quad (13)$$

and put $T = 0$ in right side of equation (12) we get the determinant value of matrix T as

$$\det(T) = \left\{ t^2 + 2t_{\mu\nu}t^{\mu\nu} \right\}^2 - 8 \left\{ \left(t^{\alpha\beta}t_{\alpha\beta} \right)^2 - 2t^{\alpha\beta}t_{\beta\gamma}t^{\gamma\delta}t_{\delta\alpha} \right\} + \frac{4}{3}it_{\alpha\beta}t_{\gamma\delta}t_{\lambda\rho}\epsilon^{\alpha\beta\gamma\delta\lambda\rho}t. \quad (14)$$

Now for trace of matrix T ,

$$\begin{aligned} \text{since } \left\{ (T-t)^2 + 2t_{\mu\nu}t^{\mu\nu} \right\}^2 &= \left\{ T^2 + t^2 - 2Tt + 2t_{\mu\nu}t^{\mu\nu} \right\}^2 \\ &= T^4 + t^4 + 4T^2t^2 + 4(t_{\mu\nu}t^{\mu\nu})^2 + 2T^2t^2 - 4T^3t + 4T^2(t_{\mu\nu}t^{\mu\nu}) \\ &\quad - 4Tt^3 + 4t^2(t_{\mu\nu}t^{\mu\nu}) - 8Tt(t_{\mu\nu}t^{\mu\nu}), \end{aligned}$$

so equation (12) becomes

$$\begin{aligned} \Delta(T) &= T^4 + t^4 + 4T^2t^2 + 4(t_{\mu\nu}t^{\mu\nu})^2 + 2T^2t^2 - 4T^3t + 4T^2(t_{\mu\nu}t^{\mu\nu}) - 4Tt^3 + 4t^2(t_{\mu\nu}t^{\mu\nu}) \\ &\quad - 8Tt(t_{\mu\nu}t^{\mu\nu}) - 8 \left\{ \left(t^{\alpha\beta}t_{\alpha\beta} \right)^2 - 2t^{\alpha\beta}t_{\beta\gamma}t^{\gamma\delta}t_{\delta\alpha} \right\} - \frac{4}{3}it_{\alpha\beta}t_{\gamma\delta}t_{\lambda\rho}\epsilon^{\alpha\beta\gamma\delta\lambda\rho} (T-t), \end{aligned}$$

differentiate above equation three times with respect to T then put $T = 0$ and divide by -6 we get trace of matrix T .

$$\frac{d^3\Delta(T)}{dT^3} = 24T - 24t, \quad \left[\frac{d^3\Delta(T)}{dT^3} \right]_{T=0} = 0 - 24t, \quad \therefore tr(T) = \frac{1}{-6} \left[\frac{d^3(\Delta(T))}{dT^3} \right]_{T=0} = 4t \quad (15)$$

equation (15) agrees with the result given in papers [1].

3. CONCLUSION:

The above reckoning helps in studying the whole paper [1] and gives a new method to find out the characteristic equation and determinant value of any 4×4 matrix in tensor form with the help of Harish-Chandra Formula for the product of Dirac-matrices.

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